A weak form of self-testing

Jędrzej Kaniewski

Center for Theoretical Physics, Polish Academy of Sciences $(\rightarrow$ Faculty of Physics, University of Warsaw)

www.jkaniewski.eu

CEQIP '19 6 June 2019

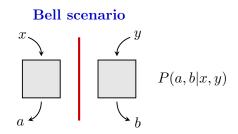


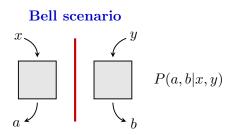
- Bell nonlocality
- Strong self-testing (CHSH)
- Weak self-testing
- Certifying randomness
- Conclusions and open questions

Outline

• Bell nonlocality

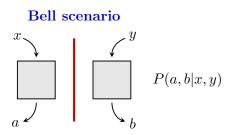
- Strong self-testing (CHSH)
- Weak self-testing
- Certifying randomness
- Conclusions and open questions





Assume that $P \in \mathcal{Q}$ is quantum

$$P(a,b|x,y) = \langle F_a^x \otimes G_b^y, \rho_{AB} \rangle.$$



Assume that $P \in \mathcal{Q}$ is quantum

$$P(a,b|x,y) = \langle F_a^x \otimes G_b^y, \rho_{AB} \rangle.$$

Definition: $P \in \mathcal{L}$ is local if

$$P(a,b|x,y) = \sum_{\lambda} p(\lambda) p_A(a|x,\lambda) p_B(b|y,\lambda).$$

Bell: $\mathcal{L} \subsetneq \mathcal{Q} \iff$ " quantum mechanics is (Bell) **nonlocal** "

Given some $P \in \mathcal{Q}$, how to show that $P \notin \mathcal{L}$?

Given some $P \in \mathcal{Q}$, how to show that $P \notin \mathcal{L}$? Real vector $C = (c_{abxy})$ define

$$\langle C, P \rangle := \sum_{abxy} c_{abxy} P(a, b | x, y)$$

and

$$\begin{split} \beta_{\mathcal{L}} &:= \max_{P \in \mathcal{L}} \left\langle C, P \right\rangle \quad \text{(local value)} \\ \beta_{\mathcal{Q}} &:= \max_{P \in \mathcal{Q}} \left\langle C, P \right\rangle \quad \text{(quantum value)} \end{split}$$

(suppose $\beta_{\mathcal{L}} < \beta_{\mathcal{Q}}$)

Given some $P \in \mathcal{Q}$, how to show that $P \notin \mathcal{L}$? Real vector $C = (c_{abxy})$ define

$$\langle C, P \rangle := \sum_{abxy} c_{abxy} P(a, b | x, y)$$

and

$$\begin{split} \beta_{\mathcal{L}} &:= \max_{P \in \mathcal{L}} \left\langle C, P \right\rangle \quad \text{(local value)} \\ \beta_{\mathcal{Q}} &:= \max_{P \in \mathcal{Q}} \left\langle C, P \right\rangle \quad \text{(quantum value)} \end{split}$$

(suppose $\beta_{\mathcal{L}} < \beta_{\mathcal{Q}}$)

Bell violation: $\langle C, P \rangle > \beta_{\mathcal{L}} \implies P \notin \mathcal{L}$

Observation: Separable states give local statistics (for all measurements)

$$\rho_{AB} = \sum_{\lambda} p_{\lambda} \sigma_{\lambda} \otimes \tau_{\lambda},$$

Observation: Separable states give local statistics (for all measurements)

$$\rho_{AB} = \sum_{\lambda} p_{\lambda} \sigma_{\lambda} \otimes \tau_{\lambda},$$

$$P(a,b|x,y) = \langle F_a^x \otimes G_b^y, \rho_{AB} \rangle = \sum_{\lambda} p_{\lambda} \cdot \underbrace{\langle F_a^x, \sigma_{\lambda} \rangle}_{p_A(a|x,\lambda)} \cdot \underbrace{\langle G_b^y, \tau_{\lambda} \rangle}_{p_B(b|y,\lambda)}.$$

Nonlocality \implies entanglement

Observation: Separable states give local statistics (for all measurements)

$$\rho_{AB} = \sum_{\lambda} p_{\lambda} \sigma_{\lambda} \otimes \tau_{\lambda},$$

$$P(a,b|x,y) = \langle F_a^x \otimes G_b^y, \rho_{AB} \rangle = \sum_{\lambda} p_{\lambda} \cdot \underbrace{\langle F_a^x, \sigma_{\lambda} \rangle}_{p_A(a|x,\lambda)} \cdot \underbrace{\langle G_b^y, \tau_{\lambda} \rangle}_{p_B(b|y,\lambda)}.$$

Nonlocality \implies entanglement

Can we make this connection more explicit/quantitative?

Outline

• Bell nonlocality

• Strong self-testing (CHSH)

- Weak self-testing
- Certifying randomness
- Conclusions and open questions

Remark 1. Bell functional vs. Bell operator

• For a Bell functional (c_{abxy}) define the Bell operator as

$$W := \sum_{abxy} c_{abxy} F_a^x \otimes G_b^y.$$

Remark 1. Bell functional vs. Bell operator

• For a Bell functional (c_{abxy}) define the Bell operator as

$$W := \sum_{abxy} c_{abxy} F_a^x \otimes G_b^y.$$

• Easy to check that

$$\langle C, P \rangle = \sum_{abxy} c_{abxy} P(a, b | x, y) = \sum_{abxy} c_{abxy} \langle F_a^x \otimes G_b^y, \rho_{AB} \rangle = \langle W, \rho_{AB} \rangle.$$

Remark 1. Bell functional vs. Bell operator

• For a Bell functional (c_{abxy}) define the Bell operator as

$$W := \sum_{abxy} c_{abxy} F_a^x \otimes G_b^y.$$

• Easy to check that

$$\langle C, P \rangle = \sum_{abxy} c_{abxy} P(a, b | x, y) = \sum_{abxy} c_{abxy} \langle F_a^x \otimes G_b^y, \rho_{AB} \rangle = \langle W, \rho_{AB} \rangle.$$

• Proving a bound on the quantum value $\beta_Q \leq c$ is equivalent to showing that

$$W \leq c \mathbb{1}$$

for all possible measurements choices.

• Measurement: resolution of 1 into positive semidefinite operators

- $\bullet\,$ Measurement: resolution of $\mathbbm{1}$ into positive semidefinite operators
- Measurements with two outcomes, i.e.

$$F_a = F_a^{\dagger}, \quad F_a \ge 0, \quad F_0 + F_1 = \mathbb{1}$$

- $\bullet\,$ Measurement: resolution of $\mathbbm{1}$ into positive semidefinite operators
- Measurements with two outcomes, i.e.

$$F_a = F_a^{\dagger}, \quad F_a \ge 0, \quad F_0 + F_1 = \mathbb{1}$$

are conveniently written as **observables**

$$A = F_0 - F_1.$$

- $\bullet\,$ Measurement: resolution of $\mathbbm{1}$ into positive semidefinite operators
- Measurements with two outcomes, i.e.

$$F_a = F_a^{\dagger}, \quad F_a \ge 0, \quad F_0 + F_1 = \mathbb{1}$$

are conveniently written as **observables**

$$A = F_0 - F_1.$$

• This mapping is one-to-one: any A such that

$$A = A^{\dagger}$$
 and $-\mathbb{1} \le A \le \mathbb{1}$

defines a valid measurement.

• The CHSH operator reads

 $W := A_0 \otimes (B_0 + B_1) + A_1 \otimes (B_0 - B_1),$

where $-\mathbb{1} \leq A_j \leq \mathbb{1}$ and $-\mathbb{1} \leq B_k \leq \mathbb{1}$.

• The CHSH operator reads

$$W := A_0 \otimes (B_0 + B_1) + A_1 \otimes (B_0 - B_1),$$

where $-\mathbb{1} \leq A_j \leq \mathbb{1}$ and $-\mathbb{1} \leq B_k \leq \mathbb{1}$.

• Well known that $\beta_{\mathcal{L}} = 2$ and $\beta_{\mathcal{Q}} = 2\sqrt{2}$.

"The maximal violation $\beta = 2\sqrt{2}$ can be achieved in an essentially unique manner"

[Tsirelson '87], [Summers and Werner '87], [Popescu and Rohrlich '92]

Proof:

• Define

$$V_0 = A_0 \otimes \mathbb{1} - \mathbb{1} \otimes \frac{B_0 + B_1}{\sqrt{2}},$$

$$V_1 = A_1 \otimes \mathbb{1} - \mathbb{1} \otimes \frac{B_0 - B_1}{\sqrt{2}}.$$

Proof:

• Define

$$V_0 = A_0 \otimes \mathbb{1} - \mathbb{1} \otimes \frac{B_0 + B_1}{\sqrt{2}},$$

$$V_1 = A_1 \otimes \mathbb{1} - \mathbb{1} \otimes \frac{B_0 - B_1}{\sqrt{2}}.$$

 \bullet Check

$$W = \frac{1}{\sqrt{2}} \left[(A_0^2 + A_1^2) \otimes \mathbb{1} + \mathbb{1} \otimes (B_0^2 + B_1^2) - (V_0^{\dagger} V_0 + V_1^{\dagger} V_1) \right].$$

Proof:

• Define

$$V_0 = A_0 \otimes \mathbb{1} - \mathbb{1} \otimes \frac{B_0 + B_1}{\sqrt{2}},$$

$$V_1 = A_1 \otimes \mathbb{1} - \mathbb{1} \otimes \frac{B_0 - B_1}{\sqrt{2}}.$$

• Check

$$W = \frac{1}{\sqrt{2}} \left[(A_0^2 + A_1^2) \otimes \mathbb{1} + \mathbb{1} \otimes (B_0^2 + B_1^2) - (V_0^{\dagger} V_0 + V_1^{\dagger} V_1) \right].$$

• Conclude that $W \leq 2\sqrt{2} \mathbb{1} \implies \operatorname{tr}(W\rho_{AB}) \leq 2\sqrt{2} = \beta_{\mathcal{Q}}$, so the SOS decomposition is tight.

$$W = \frac{1}{\sqrt{2}} \left[(A_0^2 + A_1^2) \otimes \mathbb{1} + \mathbb{1} \otimes (B_0^2 + B_1^2) - (V_0^{\dagger} V_0 + V_1^{\dagger} V_1) \right]$$

Observing $tr(W\rho_{AB}) = 2\sqrt{2}$ implies that:

$$W = \frac{1}{\sqrt{2}} \left[(A_0^2 + A_1^2) \otimes \mathbb{1} + \mathbb{1} \otimes (B_0^2 + B_1^2) - (V_0^{\dagger} V_0 + V_1^{\dagger} V_1) \right]$$

Observing $tr(W\rho_{AB}) = 2\sqrt{2}$ implies that:

• All measurements are projective on the local supports: $\operatorname{tr}(A_x^2 \rho_A) = \operatorname{tr}(B_y^2 \rho_B) = 1.$

$$W = \frac{1}{\sqrt{2}} \left[(A_0^2 + A_1^2) \otimes \mathbb{1} + \mathbb{1} \otimes (B_0^2 + B_1^2) - (V_0^{\dagger} V_0 + V_1^{\dagger} V_1) \right]$$

Observing $tr(W\rho_{AB}) = 2\sqrt{2}$ implies that:

- All measurements are projective on the local supports: $\operatorname{tr}(A_x^2 \rho_A) = \operatorname{tr}(B_y^2 \rho_B) = 1.$
- **2** Observables of Alice and Bob satisfy $V_j \rho_{AB} = 0$, e.g.

$$(A_0 \otimes \mathbb{1})\rho_{AB} = \left(\mathbb{1} \otimes \frac{B_0 + B_1}{\sqrt{2}}\right)\rho_{AB}.$$

$$W = \frac{1}{\sqrt{2}} \left[(A_0^2 + A_1^2) \otimes \mathbb{1} + \mathbb{1} \otimes (B_0^2 + B_1^2) - (V_0^{\dagger} V_0 + V_1^{\dagger} V_1) \right]$$

Observing $tr(W\rho_{AB}) = 2\sqrt{2}$ implies that:

- All measurements are projective on the local supports: $\operatorname{tr}(A_x^2 \rho_A) = \operatorname{tr}(B_y^2 \rho_B) = 1.$
- **2** Observables of Alice and Bob satisfy $V_j \rho_{AB} = 0$, e.g.

$$(A_0 \otimes \mathbb{1})\rho_{AB} = \left(\mathbb{1} \otimes \frac{B_0 + B_1}{\sqrt{2}}\right)\rho_{AB}.$$

If ρ_A and ρ_B are full-rank, then

$$A_0^2 = \mathbb{1} \implies \left(\frac{B_0 + B_1}{\sqrt{2}}\right)^2 = \mathbb{1} \implies \{B_0, B_1\} = 0.$$

• These algebraic relations determine the form of observables

$$B_0^2 = B_1^2 = 1 \quad \text{and} \quad \{B_0, B_1\} = 0 \implies \begin{array}{c} B_0 = U_B(\sigma_x \otimes 1)U_B^{\dagger} \\ B_1 = U_B(\sigma_z \otimes 1)U_B^{\dagger} \end{array}$$

.

• These algebraic relations determine the form of observables

$$B_0^2 = B_1^2 = 1 \text{ and } \{B_0, B_1\} = 0 \implies \begin{array}{c} B_0 = U_B(\sigma_x \otimes 1)U_B^{\dagger} \\ B_1 = U_B(\sigma_z \otimes 1)U_B^{\dagger} \end{array}$$

- By symmetry A_0 and A_1 have the same form.
- Construct W and determine the eigenspace corresponding to $\lambda = 2\sqrt{2}$ (essentially a two-qubit operator).

• These algebraic relations determine the form of observables

$$B_0^2 = B_1^2 = 1 \quad \text{and} \quad \{B_0, B_1\} = 0 \implies \begin{array}{c} B_0 = U_B(\sigma_x \otimes 1)U_B^{\dagger} \\ B_1 = U_B(\sigma_z \otimes 1)U_B^{\dagger} \end{array}$$

- By symmetry A_0 and A_1 have the same form.
- Construct W and determine the eigenspace corresponding to $\lambda = 2\sqrt{2}$ (essentially a two-qubit operator).

Self-testing (rigidity) statement for CHSH: if $\beta = 2\sqrt{2}$ then

$$A_0 = U_A(\sigma_x \otimes \mathbb{1})U_A^{\dagger} \qquad B_0 = U_B(\sigma_x \otimes \mathbb{1})U_B^{\dagger}$$
$$A_1 = U_A(\sigma_z \otimes \mathbb{1})U_A^{\dagger} \qquad B_1 = U_B(\sigma_z \otimes \mathbb{1})U_B^{\dagger}$$

and

 $\rho_{AB} = U(\Phi_{A'B'} \otimes \tau_{A''B''})U^{\dagger} \text{ for } U := U_A \otimes U_B$

Strategy:

- Find tight SOS decomposition.
- **2** Deduce algebraic relations between local observables.
- Obduce their exact form (up to unitaries and extra degrees of freedom).
- **(**) Construct Bell operator and find eigenspace corresponding to β_{Q} .

Strategy:

- Find tight SOS decomposition.
- **2** Deduce algebraic relations between local observables.
- Obduce their exact form (up to unitaries and extra degrees of freedom).
- **(**) Construct Bell operator and find eigenspace corresponding to β_{Q} .

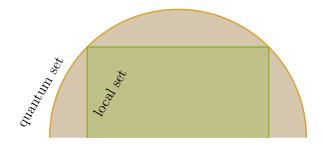


Outline

- Bell nonlocality
- Strong self-testing (CHSH)
- Weak self-testing
- Certifying randomness
- Conclusions and open questions

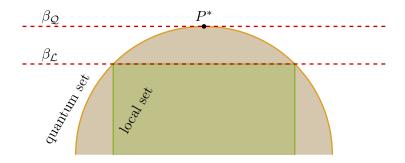
A weak form of self-testing

Immediate consequence of self-testing: the maximal violation is achieved by a unique probability point



What about Bell functionals that do not have a unique maximiser?

Immediate consequence of self-testing: the maximal violation is achieved by a unique probability point

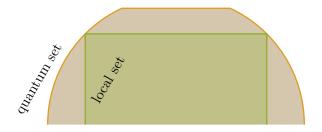


What about Bell functionals that do not have a unique maximiser?

• In a scenario with 3 inputs and 2 outputs per party consider $W = A_0 \otimes (B_0 + B_1 + B_2) + A_1 \otimes (B_0 + B_1 - B_2) + A_2 \otimes (B_0 - B_1).$

(the correlation part of the infamous I_{3322} functional)

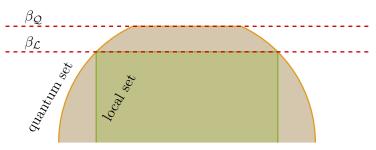
• Easy to show that $\beta_{\mathcal{L}} = 4$ and $\beta_{\mathcal{Q}} = 5$, but the maximal violation is achieved by multiple probability points.



• In a scenario with 3 inputs and 2 outputs per party consider $W = A_0 \otimes (B_0 + B_1 + B_2) + A_1 \otimes (B_0 + B_1 - B_2) + A_2 \otimes (B_0 - B_1).$

(the correlation part of the infamous I_{3322} functional)

• Easy to show that $\beta_{\mathcal{L}} = 4$ and $\beta_{\mathcal{Q}} = 5$, but the maximal violation is achieved by multiple probability points.



no unique maximiser \implies no rigidity statement can we still have some **weak form of self-testing**?

• We already have a tight SOS decomposition

$$2W = (2A_0^2 + 2A_1^2 + A_2^2) \otimes \mathbb{1} + \mathbb{1} \otimes (2B_0^2 + 2B_1^2 + B_2^2) - \sum_{j=0}^2 V_j^{\dagger} V_j,$$

where

$$V_0 = (A_0 + A_1) \otimes \mathbb{1} - \mathbb{1} \otimes (B_0 + B_1),$$

$$V_1 = (A_0 - A_1) \otimes \mathbb{1} - \mathbb{1} \otimes B_2,$$

$$V_2 = A_2 \otimes \mathbb{1} - \mathbb{1} \otimes (B_0 - B_1).$$

• We already have a tight SOS decomposition

$$2W = (2A_0^2 + 2A_1^2 + A_2^2) \otimes \mathbb{1} + \mathbb{1} \otimes (2B_0^2 + 2B_1^2 + B_2^2) - \sum_{j=0}^2 V_j^{\dagger} V_j,$$

where

$$V_0 = (A_0 + A_1) \otimes \mathbb{1} - \mathbb{1} \otimes (B_0 + B_1),$$

$$V_1 = (A_0 - A_1) \otimes \mathbb{1} - \mathbb{1} \otimes B_2,$$

$$V_2 = A_2 \otimes \mathbb{1} - \mathbb{1} \otimes (B_0 - B_1).$$

• Observing $\beta = 5$ implies

$$\operatorname{tr}(A_j^2 \rho_A) = \operatorname{tr}(B_j^2 \rho_B) = 1,$$
$$V_j \rho_{AB} = 0$$

for j = 0, 1, 2.

• Let us derive an explicit form of the observables. Rewriting $V_1\rho_{AB}=0$ gives

$$[(A_0 - A_1) \otimes \mathbb{1}]\rho_{AB} = (\mathbb{1} \otimes B_2)\rho_{AB},$$

which combined with the full-rank assumption leads to

$$(A_0 - A_1)^2 = \mathbb{1}.$$

• Let us derive an explicit form of the observables. Rewriting $V_1\rho_{AB}=0$ gives

$$[(A_0 - A_1) \otimes \mathbb{1}]\rho_{AB} = (\mathbb{1} \otimes B_2)\rho_{AB},$$

which combined with the full-rank assumption leads to

$$(A_0 - A_1)^2 = \mathbb{1}.$$

• Together with projectivity this gives

$$\begin{aligned} \mathfrak{H}_A &\equiv \mathbb{C}^2 \otimes \mathbb{C}^{d_A}, \\ A_0 &= \left(\cos\frac{\pi}{6}\,\mathsf{X} + \sin\frac{\pi}{6}\,\mathsf{Z}\right) \otimes \mathbb{1}, \\ A_1 &= \left(\cos\frac{\pi}{6}\,\mathsf{X} - \sin\frac{\pi}{6}\,\mathsf{Z}\right) \otimes \mathbb{1}. \end{aligned}$$

(up to a choice of basis)

• Combining
$$V_0\rho_{AB} = 0$$
 and $V_2\rho_{AB} = 0$ gives

$$\frac{1}{2} [(A_0 + A_1 + A_2) \otimes \mathbb{1}] \rho_{AB} = (\mathbb{1} \otimes B_2)\rho_{AB}$$
and finally
$$(A_0 + A_1 + A_2)^2 = 4 \mathbb{1}.$$

• Combining
$$V_0\rho_{AB} = 0$$
 and $V_2\rho_{AB} = 0$ gives

$$\frac{1}{2} [(A_0 + A_1 + A_2) \otimes \mathbb{1}] \rho_{AB} = (\mathbb{1} \otimes B_2)\rho_{AB}$$
and finally
$$(A_0 + A_1 + A_2)^2 = 4 \mathbb{1}.$$

• Plugging in the previously derived characterisation of A_0, A_1 gives $(2\cos\frac{\pi}{6} X \otimes \mathbb{1} + A_2)^2 = 4 \mathbb{1}.$

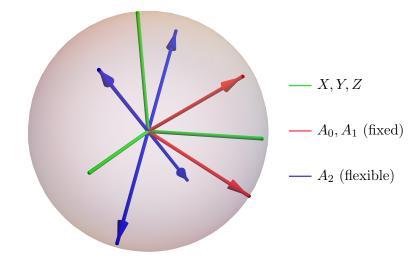
• Combining
$$V_0\rho_{AB} = 0$$
 and $V_2\rho_{AB} = 0$ gives

$$\frac{1}{2} [(A_0 + A_1 + A_2) \otimes \mathbb{1}] \rho_{AB} = (\mathbb{1} \otimes B_2)\rho_{AB}$$
and finally
$$(A_0 + A_1 + A_2)^2 = 4 \mathbb{1}.$$

- Plugging in the previously derived characterisation of A_0, A_1 gives $(2\cos\frac{\pi}{6} X \otimes \mathbb{1} + A_2)^2 = 4 \mathbb{1}.$
- Simple algebra yields

$$A_2 = \sum_{j=1}^{d_A} \left(\cos u_j \,\mathsf{Y} + \sin u_j \,\mathsf{Z} \right) \otimes |e_j\rangle \langle e_j|,$$

where $u_j \in [0, 2\pi)$ and $\{|e_j\rangle\}_{j=1}^{d_A}$ is an orthonormal basis on \mathbb{C}^{d_A} .



• By symmetry the same characterisation holds for the observables of Bob. The Bell operator reads

$$W = \sum_{j=1}^{d_A} \sum_{k=1}^{d_B} R(u_j, v_k) \otimes |e_j\rangle \langle e_j| \otimes |f_k\rangle \langle f_k|,$$

where $R(u_j, v_k)$ is a two-qubit operator.

• By symmetry the same characterisation holds for the observables of Bob. The Bell operator reads

$$W = \sum_{j=1}^{d_A} \sum_{k=1}^{d_B} R(u_j, v_k) \otimes |e_j\rangle \langle e_j| \otimes |f_k\rangle \langle f_k|,$$

where $R(u_j, v_k)$ is a two-qubit operator.

• Diagonalising $R(u_j, v_k)$ shows that $\lambda = 5$ belong to the spectrum iff $u_j = v_k$.

• By symmetry the same characterisation holds for the observables of Bob. The Bell operator reads

$$W = \sum_{j=1}^{d_A} \sum_{k=1}^{d_B} R(u_j, v_k) \otimes |e_j\rangle \langle e_j| \otimes |f_k\rangle \langle f_k|,$$

where $R(u_j, v_k)$ is a two-qubit operator.

- Diagonalising $R(u_j, v_k)$ shows that $\lambda = 5$ belong to the spectrum iff $u_j = v_k$.
- The corresponding eigenvector is the standard maximally entangled state $|\Phi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$ (smart parameterisation)

• By symmetry the same characterisation holds for the observables of Bob. The Bell operator reads

$$W = \sum_{j=1}^{d_A} \sum_{k=1}^{d_B} R(u_j, v_k) \otimes |e_j\rangle \langle e_j| \otimes |f_k\rangle \langle f_k|,$$

where $R(u_j, v_k)$ is a two-qubit operator.

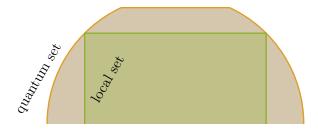
- Diagonalising $R(u_j, v_k)$ shows that $\lambda = 5$ belong to the spectrum iff $u_j = v_k$.
- The corresponding eigenvector is the standard maximally entangled state $|\Phi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$ (smart parameterisation)
- Finally, the global state must be of the form

$$\rho_{AB} = \Phi^+_{A'B'} \otimes \sigma_{A''B''},$$

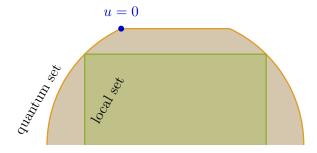
where $\sigma_{A''B''}$ is a normalised state satisfying

$$\operatorname{tr}\left(\sigma_{A''B''}|e_{j}\rangle\langle e_{j}|\otimes|f_{k}\rangle\langle f_{k}|\right)=0 \quad \text{if} \quad u_{j}\neq v_{k}.$$

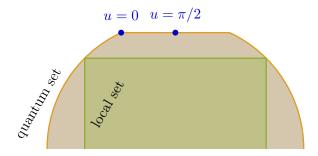
Conclusion: there is a family of optimal strategies parametrised by $u \in [0, 2\pi)$ and every optimal strategy is just a convex combination of those.



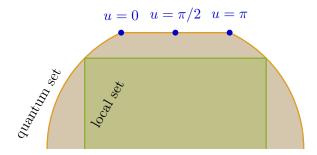
Conclusion: there is a family of optimal strategies parametrised by $u \in [0, 2\pi)$ and every optimal strategy is just a convex combination of those.



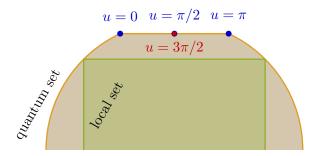
Conclusion: there is a family of optimal strategies parametrised by $u \in [0, 2\pi)$ and every optimal strategy is just a convex combination of those.



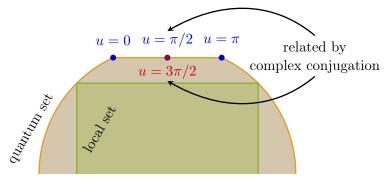
Conclusion: there is a family of optimal strategies parametrised by $u \in [0, 2\pi)$ and every optimal strategy is just a convex combination of those.



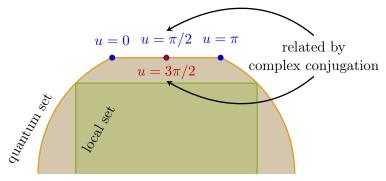
Conclusion: there is a family of optimal strategies parametrised by $u \in [0, 2\pi)$ and every optimal strategy is just a convex combination of those.



Conclusion: there is a family of optimal strategies parametrised by $u \in [0, 2\pi)$ and every optimal strategy is just a convex combination of those.



Conclusion: there is a family of optimal strategies parametrised by $u \in [0, 2\pi)$ and every optimal strategy is just a convex combination of those.



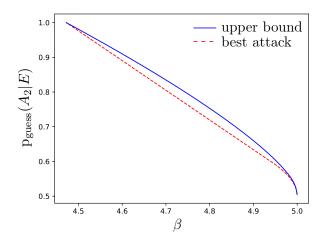
- 1. The face is a line segment.
- 2. The endpoints are self-tests in the usual strong sense.

Outline

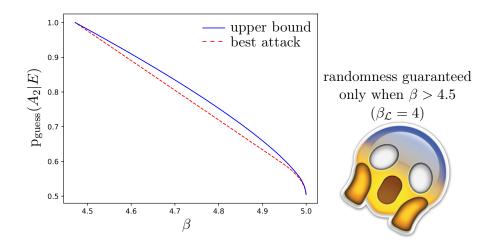
- Bell nonlocality
- Strong self-testing (CHSH)
- Weak self-testing
- Certifying randomness
- Conclusions and open questions

Setup: Alice and Bob observe Bell violation, Eve is trying to guess the outcome of Alice for a specific setting. We start with A_2 :

Setup: Alice and Bob observe Bell violation, Eve is trying to guess the outcome of Alice for a specific setting. We start with A_2 :

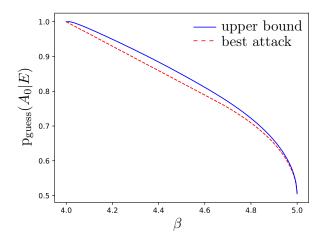


Setup: Alice and Bob observe Bell violation, Eve is trying to guess the outcome of Alice for a specific setting. We start with A_2 :

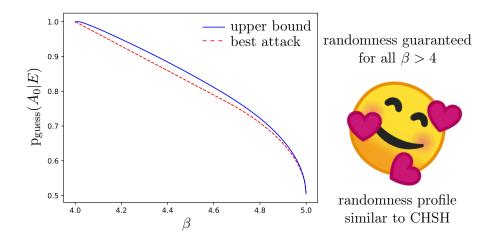


Luckily A_0 and A_1 are **much better**!

Luckily A_0 and A_1 are **much better**!



Luckily A_0 and A_1 are **much better**!



Outline

- Bell nonlocality
- Strong self-testing (CHSH)
- Weak self-testing
- Certifying randomness
- Conclusions and open questions

- A new weak form of self-testing: the maximal violation certifies the state, but not the measurements.
- Nevertheless, the randomness certification power is not significantly affected.

- A new weak form of self-testing: the maximal violation certifies the state, but not the measurements.
- Nevertheless, the randomness certification power is not significantly affected.

Open questions:

- Is the self-testing robust? How to construct an extraction channel which depends on 3 observables (instead of the usual 2)?
- Can we find a scenario in which the non-rigidty has a significant impact, e.g. for randomness certification?
- Can we find a bipartite Bell inequality which can be maximally violated by distinct states? (exists for 3 parties)

- A new weak form of self-testing: the maximal violation certifies the state, but not the measurements.
- Nevertheless, the randomness certification power is not significantly affected.

Open questions:

- Is the self-testing robust? How to construct an extraction channel which depends on 3 observables (instead of the usual 2)?
- Can we find a scenario in which the non-rigidty has a significant impact, e.g. for randomness certification?
- Can we find a bipartite Bell inequality which can be maximally violated by distinct states? (exists for 3 parties)

see also related independent work by Jebarathinam et al. arXiv:1905.09867 (based on lifting)

- A new weak form of self-testing: the maximal violation certifies the state, but not the measurements.
- Nevertheless, the randomness certification power is not significantly affected.

Open questions:

- Is the self-testing robust? How to construct an extraction channel which depends on 3 observables (instead of the usual 2)?
- Can we find a scenario in which the non-rigidty has a significant impact, e.g. for randomness certification?
- Can we find a bipartite Bell inequality which can be maximally violated by distinct states? (exists for 3 parties)

see also related independent work by Jebarathinam et al. arXiv:1905.09867 (based on lifting)

Thank you for your attention!