Analytic and nearly optimal self-testing bounds for the Clauser-Horne-Shimony-Holt and Mermin inequalities Phys. Rev. Lett. **117**, 070402 (2016) [arXiv:1604.08176]

Jed Kaniewski

QMATH, Department of Mathematical Sciences

University of Copenhagen, Universitetsparken 5, 2100 Copenhagen Ø, Denmark

17 June 2016

CEQIP '16





http://qmath.ku.dk/

- What is self-testing?
- Previous results and new findings
- Self-testing from operator inequalities
- Two examples: the CHSH and Mermin₃ inequalities
- Summary and future work

Outline

- What is self-testing?
- Previous results and new findings
- Self-testing from operator inequalities
- Two examples: the CHSH and Mermin₃ inequalities
- Summary and future work

Bell scenario



Bell scenario



Def.: Pr[a, b|x, y] is local if

$$\Pr[a, b|x, y] = \sum_{\lambda} p(\lambda) p(a|x, \lambda) p(b|y, \lambda).$$

Otherwise \implies nonlocal or it violates (some) Bell inequality

Obs.: Separable states give local statistics (for all measurements)

$$\rho_{AB} = \sum_{\lambda} p_{\lambda} \alpha_{\lambda} \otimes \beta_{\lambda},$$

$$\Pr[a, b|x, y] = \operatorname{tr}\left[(P_a^x \otimes Q_b^y)\rho_{AB}\right] = \sum_{\lambda} p_{\lambda} \cdot \underbrace{\operatorname{tr}(P_a^x \alpha_{\lambda})}_{p(a|x,\lambda)} \cdot \underbrace{\operatorname{tr}(Q_b^y \beta_{\lambda})}_{p(b|y,\lambda)}.$$

 ρ_{AB} is separable \implies statistics are local $\Pr[a, b|x, y]$ is nonlocal $\implies \rho_{AB}$ is entangled





Google-søgning Jeg prøver lykken

Given $\Pr[a, b|x, y] = \operatorname{tr}\left[(P_a^x \otimes Q_b^y)\rho_{AB}\right]$

deduce properties of ρ_{AB} , $\{P_a^x\}$, $\{Q_b^y\}$

Given $\Pr[a, b|x, y] = \operatorname{tr}\left[(P_a^x \otimes Q_b^y)\rho_{AB}\right]$

deduce properties of ρ_{AB} , $\{P_a^x\}$, $\{Q_b^y\}$ We don't assume that ρ_{AB} is pure and it's **important**! (ask me if you want to know more)

Given $\Pr[a, b|x, y] = \operatorname{tr}\left[(P_a^x \otimes Q_b^y)\rho_{AB}\right]$

deduce properties of ρ_{AB} , $\{P_a^x\}$, $\{Q_b^y\}$ We don't assume that ρ_{AB} is pure and it's **important**! (ask me if you want to know more)

often only promised some Bell violation

$$\sum_{abxy} c_{ab}^{xy} \Pr[a, b | x, y] = \beta$$

Example: the CHSH inequality [Popescu, Rohrlich '92]

$$\beta_{\text{CHSH}} := \sum_{abxy} (-1)^{a+b+xy} \Pr[a, b|x, y] \text{ for } a, b, x, y \in \{0, 1\}$$

 $\beta_{\text{CHSH}} = 2\sqrt{2} \text{ (max)} \implies \rho_{AB} \simeq \Phi_{AB} \text{ for } |\Phi_{AB}\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle).$

Example: the CHSH inequality [Popescu, Rohrlich '92]

$$\beta_{\text{CHSH}} := \sum_{abxy} (-1)^{a+b+xy} \Pr[a, b|x, y] \text{ for } a, b, x, y \in \{0, 1\}$$

 $\beta_{\text{CHSH}} = 2\sqrt{2} \text{ (max)} \implies \rho_{AB} \simeq \Phi_{AB} \text{ for } |\Phi_{AB}\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle).$

$$\rho_{AB} = \Phi_{AB}$$

Inherent limitations

Example: the CHSH inequality [Popescu, Rohrlich '92]

$$\beta_{\text{CHSH}} := \sum_{abxy} (-1)^{a+b+xy} \Pr[a, b|x, y] \text{ for } a, b, x, y \in \{0, 1\}$$

 $\beta_{\text{CHSH}} = 2\sqrt{2} \text{ (max)} \implies \rho_{AB} \simeq \Phi_{AB} \text{ for } |\Phi_{AB}\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle).$

$$\rho_{AB} = \Phi_{AB} \otimes \tau_{A'B'}$$

Inherent limitations

• cannot see auxiliary systems (measurements act trivially)

Example: the CHSH inequality [Popescu, Rohrlich '92]

$$\beta_{\text{CHSH}} := \sum_{abxy} (-1)^{a+b+xy} \Pr[a, b|x, y] \text{ for } a, b, x, y \in \{0, 1\}$$

 $\beta_{\text{CHSH}} = 2\sqrt{2} \text{ (max)} \implies \rho_{AB} \simeq \Phi_{AB} \text{ for } |\Phi_{AB}\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle).$

$$\rho_{AB} = U(\Phi_{AB} \otimes \tau_{A'B'})U^{\dagger} \quad \text{for} \quad U = U_{AA'} \otimes U_{BB'}$$

Inherent limitations

- cannot see auxiliary systems (measurements act trivially)
- cannot see local unitaries

Example: the CHSH inequality [Popescu, Rohrlich '92]

$$\beta_{\text{CHSH}} := \sum_{abxy} (-1)^{a+b+xy} \Pr[a, b|x, y] \text{ for } a, b, x, y \in \{0, 1\}$$

 $\beta_{\text{CHSH}} = 2\sqrt{2} \text{ (max)} \implies \rho_{AB} \simeq \Phi_{AB} \text{ for } |\Phi_{AB}\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle).$

$$\rho_{AB} = U(\Phi_{AB} \otimes \tau_{A'B'})U^{\dagger} \quad \text{for} \quad U = U_{AA'} \otimes U_{BB'}$$

Inherent limitations

- cannot see auxiliary systems (measurements act trivially)
- cannot see local unitaries

$$\sum_{abxy} c_{ab}^{xy} \Pr[a, b|x, y] = \beta$$



state certification









What is **experimentally-relevant**?

The CHSH inequality: $\beta_C = 2$ and $\beta_Q = 2\sqrt{2}$

What is **experimentally-relevant**?

The CHSH inequality: $\beta_C = 2$ and $\beta_Q = 2\sqrt{2}$

Non-trivial bounds for...

[Bardyn et al. '09]: $\beta \geq 1 + \sqrt{2} \approx 2.41$

[McKague et al. '12]: $\beta \geq \beta_Q - 2 \cdot 10^{-5}$

What is **experimentally-relevant**?

The CHSH inequality: $\beta_C = 2$ and $\beta_Q = 2\sqrt{2}$

Non-trivial bounds for...

[Bardyn et al. '09]: $\beta \geq 1 + \sqrt{2} \approx 2.41$

[McKague et al. '12]: $\beta \geq \beta_Q - 2 \cdot 10^{-5}$

What is **experimentally-relevant**?

The CHSH inequality: $\beta_C = 2$ and $\beta_Q = 2\sqrt{2}$

Non-trivial bounds for...

[Bardyn et al. '09]: $\beta \geq 1 + \sqrt{2} \approx 2.41$

[McKague et al. '12]:
$$\beta \ge \beta_Q - 2 \cdot 10^{-5}$$

The loophole-free Bell experiment from Delft

$$\beta = 2.4 \pm 0.2$$

What is **experimentally-relevant**?

The CHSH inequality: $\beta_C = 2$ and $\beta_Q = 2\sqrt{2}$

Non-trivial bounds for...

[Bardyn et al. '09]: $\beta \geq 1 + \sqrt{2} \approx 2.41$

[McKague et al. '12]:
$$\beta \ge \beta_Q - 2 \cdot 10^{-5}$$

The loophole-free Bell experiment from Delft

$$\beta = 2.4 \pm 0.2$$

• Previous results and new findings

- Self-testing from operator inequalities
- \bullet Two examples: the CHSH and Mermin_3 inequalities
- Summary and future work

Previous results

Self-testable

- the singlet [McKague et al. '12]
- graph states [McKague '14]
- high-dimensional maximally entangled state [Slofstra '11, Yang, Navascués '13, McKague '16, Salavrakos et al. '16 + ...]
- non-maximally entangled states of 2 qubits [Bamps, Pironio '15]

Only for almost perfect statistics ($\varepsilon \approx 10^{-4}$).

Previous results

Self-testable

- the singlet [McKague et al. '12]
- graph states [McKague '14]
- high-dimensional maximally entangled state [Slofstra '11, Yang, Navascués '13, McKague '16, Salavrakos et al. '16 + ...]
- non-maximally entangled states of 2 qubits [Bamps, Pironio '15]

Only for almost perfect statistics ($\varepsilon \approx 10^{-4}$).

Experimentally-relevant robustness

- a single analytic result for the singlet-CHSH case [Bardyn et al. '09]
- swap trick: a numerical method, versatile but computationally expensive (so far up to 4 qubits or 2 qutrits) [Yang et al. '14, Bancal et al. '15]

[see arXiv:1604.08176 for references]

New approach for analytic self-testing bounds

- \bullet improvement for the CHSH and Mermin_3
- Mermin₃ is actually **tight** (!)
- $\bullet\,$ self-testing problem \leadsto operator inequalities

- What is self-testing?
- Previous results and new findings
- Self-testing from operator inequalities
- Two examples: the CHSH and Mermin₃ inequalities
- Summary and future work

Extractability of $\Psi_{A'B'}$ from ρ_{AB}

$$\Xi(\rho_{AB} \to \Psi_{A'B'}) := \max_{\Lambda_A, \Lambda_B} F((\Lambda_A \otimes \Lambda_B)(\rho_{AB}), \Psi_{A'B'})$$

ocal extraction channels fidelity

Obs1: $\Xi(\rho_{AB} \to \Psi_{A'B'}) = 1 \iff \rho_{AB} = V(\Psi_{A'B'} \otimes \sigma_{A''B''})V^{\dagger}$

for $V = V_{A'A'' \to A} \otimes V_{B'B'' \to B}$

Idea: measurement operators → extraction channels!
Analytical bound of [Bardyn et al.] in 2 steps
[1] Solve the problem for 2 qbits
(local measurements determine a local unitary correction)
[2] Use Jordan's lemma to argue that it holds in arbitrary dimension

Idea: measurement operators → extraction channels!
Analytical bound of [Bardyn et al.] in 2 steps
[1] Solve the problem for 2 qbits
(local measurements determine a local unitary correction)
[2] Use Jordan's lemma to argue that it holds in arbitrary dimension

Refined approach: assume

$$\Lambda_A := \Lambda_A(\{P_a^x\})$$
 and $\Lambda_B := \Lambda_B(\{Q_b^y\}).$

Refined approach: assume

$$\Lambda_A := \Lambda_A(\{P_a^x\}) \quad \text{and} \quad \Lambda_B := \Lambda_B(\{Q_b^y\}).$$

for $\Psi_{A'B'}$ pure

$$F((\Lambda_A \otimes \Lambda_B)(\rho_{AB}), \Psi_{A'B'}) = \langle (\Lambda_A \otimes \Lambda_B)(\rho_{AB}), \Psi_{A'B'} \rangle$$
$$= \langle \rho_{AB}, (\Lambda_A^{\dagger} \otimes \Lambda_B^{\dagger})(\Psi_{A'B'}) \rangle = \operatorname{tr}(K\rho_{AB})$$

 for

$$K = (\Lambda^{\dagger}_A \otimes \Lambda^{\dagger}_B)(\Psi_{A'B'})$$

important: K depends only on $\Psi_{A'B'}$, $\{P_a^x\}$, $\{Q_b^y\}$, not on ρ_{AB} !

Refined approach: assume

$$\Lambda_A := \Lambda_A(\{P_a^x\}) \quad \text{and} \quad \Lambda_B := \Lambda_B(\{Q_b^y\}).$$

for $\Psi_{A'B'}$ pure

$$F((\Lambda_A \otimes \Lambda_B)(\rho_{AB}), \Psi_{A'B'}) = \langle (\Lambda_A \otimes \Lambda_B)(\rho_{AB}), \Psi_{A'B'} \rangle$$
$$= \langle \rho_{AB}, (\Lambda_A^{\dagger} \otimes \Lambda_B^{\dagger})(\Psi_{A'B'}) \rangle = \operatorname{tr}(K\rho_{AB})$$

for

$$K=(\Lambda^\dagger_A\otimes\Lambda^\dagger_B)(\Psi_{A'B'})$$

important: K depends only on $\Psi_{A'B'}$, $\{P_a^x\}$, $\{Q_b^y\}$, not on ρ_{AB} ! ... just like the **Bell operator**

$$W = \sum_{abxy} c_{ab}^{xy} P_a^x \otimes Q_b^y.$$

Forget the input state ρ_{AB} ! Want $s, \mu \in \mathbb{R}$ such that

 $K \geq sW + \mu \mathbb{1}$

holds for all possible measurement operators

Forget the input state ρ_{AB} ! Want $s, \mu \in \mathbb{R}$ such that

 $K \geq sW + \mu \mathbb{1}$

holds for all possible measurement operators

Challenging!... but if works then

$$\operatorname{tr}(K\rho_{AB}) \ge s \operatorname{tr}(W\rho_{AB}) + \mu \operatorname{tr}(\rho_{AB})$$

equivalent to

$$F((\Lambda_A \otimes \Lambda_B)(\rho_{AB}), \Psi_{A'B'}) \ge s\beta + \mu$$

precisely a (linear) self-testing statement!

Main technical challenge: find channels and s, μ such that

 $K \geq sW + \mu \mathbb{1}$

holds for all possible measurement operators

Main technical challenge: find channels and s, μ such that

 $K \geq sW + \mu \mathbb{1}$

holds for all possible measurement operators

Jordan's lemma: any two binary, projective measurements can be simultaneously block-diagonalised into 2×2 blocks (at most)

each block parametrised by an angle $a \in [0, \pi/2]$ (up to unitary)

this becomes tractable: 1-parameter per party

- What is self-testing?
- Previous results and new findings
- Self-testing from operator inequalities
- Two examples: the CHSH and Mermin₃ inequalities
- Summary and future work

• Extraction channels: angle-dependent dephasing

• Extraction channels: angle-dependent dephasing

• Extraction channels: angle-dependent dephasing

 \blacksquare Find suitable s,μ (numerics): $s=(4+5\sqrt{2})/16$ and $\mu=-(1+2\sqrt{2})/4$

In Prove

$$K(a,b) \geq sW(a,b) + \mu \mathbb{1}$$

for all $a, b \in [0, \pi/2]$ (2-parameter family of 4×4 matrices)

- Same extraction channels
- **2** Find suitable s, μ (numerics): $s = (2 + \sqrt{2})/8$ and $\mu = -1/\sqrt{2}$
- In Prove

$$K(a,b,c) \geq sW(a,b,c) + \mu \mathbb{1}$$

for all $a, b, c \in [0, \pi/2]$ (3-parameter family of 8×8 matrices)

$Mermin_3$ self-testing: proof in 4 steps

$Mermin_3$ self-testing: proof in 4 steps

- What is self-testing?
- Previous results and new findings
- Self-testing from operator inequalities
- Two examples: the CHSH and Mermin₃ inequalities
- Summary and future work

Summary

- self-testing from operator inequalities
- \bullet improvements for the CHSH and $Mermin_3$ inequalities
- first provably tight self-testing statement

Summary

- self-testing from operator inequalities
- \bullet improvements for the CHSH and Mermin $_3$ inequalities
- first provably tight self-testing statement

Future work

- Mermin_n $\stackrel{?}{\Longrightarrow}$ **GHZ**_n state (preliminary numerics)
- tilted CHSH $\stackrel{?}{\Longrightarrow}$ non-maximally entangled 2-qubit states [project in progress with Tim Coopmans and Christian Schaffner]
- beyond **Jordan's lemma**?
- apply this approach to **steering**?

Yes, Pooh, quantum mechanics is very strange and nobody really understands it but let's talk about it some other day...

So you can really certify quantum states without trusting the devices at all?