

\mathcal{H} - Hilbert space
 ρ - pure state
 $g, \sigma \in \mathcal{L}(\mathcal{H})$ mixed state
 measurement on \mathcal{H} with n
 $\{F_j\}_{j=1}^n$ $F_j \in \mathcal{L}(\mathcal{H})$ outcomes
 $F_j \geq 0, \sum_{j=1}^n F_j = \mathbb{1}$
 $\{K_j\}_{j=1}^n, \sum_{j=1}^n K_j K_j^\dagger = \mathbb{1}$ $F_j = K_j^\dagger K_j$
 $g \mapsto \sum_{j=1}^n K_j g K_j^\dagger \otimes |j\rangle\langle j|$
↑ ↑
perturbations classical
state outcome

Historical introduction



Alice measures in $\{|0\rangle, |1\rangle\}$
 \Rightarrow Bob's system collapsed to
 be either in $|0\rangle$ or $|1\rangle$
 $|0\rangle$ or $|1\rangle$
 \Rightarrow Alice can perfectly predict
 result of Bob's $\{|0\rangle, |1\rangle\}$
 \Rightarrow there should exist signals
classical communication

Not consistent with uncertainty
 principle
 EPR: QM formulation not complete
 realism: QM does not satisfy
 local realism
 realism: objects have properties
 (regardless of measurement)
 local: those properties should be
 affected by events far away
 with no time delay

Bell experiment



$x, y \in \{0, 1\} = \{A, B, \dots, k\}$
 $a, b \in \{0, 1\}$
 - boxes can't communicate
 - can repeat many times,
 boxes are reliable
 $P(a, b | x, y)$ - well-defined, can
 be calculated to arbitrary precision
 $P = \{P(a, b | x, y)\}_{x, y} \in \mathbb{R}^{n \times n}$
 probability joint

What happens if λ for the physical theory?

Classical theories



λ - shared random variable
 $r_A(a | x, \lambda) : \{x\} \times \{k\} \times \mathcal{L} \rightarrow \mathbb{R}_+$
 $\sum_a r_A(a | x, \lambda) = 1$
 $r_B(b | y, \lambda)$

$P(a, b | x, y) = \sum_{\lambda \in \mathcal{L}} q(\lambda) \cdot r_A(a | x, \lambda) \cdot r_B(b | y, \lambda)$

P is local mathematical definition of local realism
 P admits a local hidden variable (LHV) model
 \mathcal{L} - set of all P 's of the form above
 $|\mathcal{L}| = 1$
 $P(a, b | x, y) = r_A(a | x) \cdot r_B(b | y)$
uncorrelated
 \mathcal{L} = convex combinations of uncorrelated distributions (i.e. separable states)